## Lecture 10 (Unit-II) ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Basics of electromagnetic theory
$>$ Equation of continuity

## INTRODUCTION

$>$ Initially, we studied electricity (when charge is at rest) and magnetism (when charge is in uniform motion).
$>$ In 1820, Oersted $\longrightarrow$ magnetic field can be produced by electric current.
$>$ Faraday $\longrightarrow$ phenomenon of electromagnetic induction and showed that electric current can be produced by a time varying magnetic field.
$>$ In 1864, Maxwell unified both electric and magnetic field $\longrightarrow$ An accelerated charge particle generates electromagnetic waves.
$>$ The coupled oscillating electric and magnetic field that moves with the speed of light and exhibit wave behavior is called Electromagnetic Wave. The microwaves, infrared rays, ultraviolet rays, X-rays and $\gamma$ - rays are few examples of electromagnetic waves.
> NOTE: The electric field and magnetic field of an electromagnetic wave are perpendicular to each other and also perpendicular to the direction of propagation of EM waves.


## Some important tools

## $>$ Dell Operator

$\vec{\nabla}=\frac{\partial}{\partial \mathrm{x}} \hat{\imath}+\frac{\partial}{\partial \mathrm{y}} \hat{\jmath}+\frac{\partial}{\partial \mathrm{z}} \hat{k} \quad$ (Vector quantity)
> Laplacian Operator
$\nabla^{2}=\vec{\nabla} \cdot \vec{\nabla}=\frac{\partial^{2}}{\partial \mathrm{x}^{2}}+\frac{\partial^{2}}{\partial \mathrm{y}^{2}}+\frac{\partial^{2}}{\partial \mathrm{z}^{2}}$ (Scalar quantity)
$>$ Gradiant of $\phi=\operatorname{grad} \phi=\overrightarrow{\boldsymbol{\nabla}} \phi$
$>$ Divergence of $\overrightarrow{\boldsymbol{A}}=\operatorname{div} \overrightarrow{\boldsymbol{A}}=\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{A}}$
$>\operatorname{Curlof} \overrightarrow{\boldsymbol{A}}=\operatorname{Curl} \overrightarrow{\boldsymbol{A}}=\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{A}}$

## Some important tools

## State Gauss divergence and Stoke's theorem.

Gauss Divergence theorem - Surface integral of a vector over a closed surface area is equal to volume integral of divergence of the same vector over the volume enclosed by that surface area, i.e.

$$
\oint \overrightarrow{\boldsymbol{A}} \cdot \overrightarrow{\boldsymbol{d S}}=\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{A}}) d V
$$

## Some important tools

Stoke's theorem - Line integral of a vector over a closed loop is equal to surface integral of curl of the same vector over the surface area enclosed by that loop, i.e.

$$
\oint \vec{A} \cdot \overrightarrow{d l}=\int(\vec{\nabla} \times \vec{A}) \cdot \overrightarrow{d S}
$$

What is the equation of continuity? Obtain the required expression for it. Also, give its physical significance. [2016-17, 2018-19]

## CONTINUITY EQUATION

Electric current through a closed surface area is given by

$$
\begin{equation*}
I=\oint \overrightarrow{\boldsymbol{J}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}} \tag{1}
\end{equation*}
$$

where $\vec{J}$ is current density
Consider, a closed surface $S$ enclosing a volume $V$. If $\rho$ is the volume density inside the volume, then

$$
\begin{equation*}
I=-\frac{d q}{d t}=-\frac{d}{d t} \int \rho d V \tag{2}
\end{equation*}
$$

-ve sign indicate that charge decreases inside the volume with time.

If there is neither source nor sink inside the volume, then according to the law of conservation of charge value of electric current given by equations (1) and
(2) will be same.Therefore
$\oint \overrightarrow{\boldsymbol{J}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=-\frac{d}{d t} \int \rho d V$
Applying Gauss divergence theorem on left hand side we get

$$
\int(\vec{\nabla} \cdot \vec{J}) d V=-\frac{d}{d t} \int \rho d V
$$

$$
\begin{aligned}
& \int(\vec{\nabla} \cdot \vec{J}) d V=-\int \frac{\partial \rho}{\partial t} d V \\
& \int\left(\vec{\nabla} \cdot \vec{J}+\frac{\partial \rho}{\partial t}\right) d V=0
\end{aligned}
$$

For any arbitrary volume, the integral must be zero. Thus,

$$
\vec{\nabla} \cdot \vec{J}+\frac{\partial \rho}{\partial t}=0
$$

This is called equation of continuity.

## CONTINUITY EQUATION

## > Physical significance:

Law of conservation of charge: It states that the total current flowing out of some volume must be equal to the rate of decrease of charge within the volume, assuming that charge can neither be created nor be destroyed.
$>$ For steady state

$$
\begin{array}{cl}
\frac{\partial \rho}{\partial t}=0 & \\
\therefore & \vec{\nabla} \cdot \vec{J}=0
\end{array}
$$

## Lecture 11 ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Displacement current
> Modification of Ampere's circuital law

Why Maxwell proposed that Ampere's law require modification? [2018-19]

The idea of modification in Ampere's law arises in connection with capacitors with no medium between them.
Let us consider the circuit showing the process of charging of capacitor. Consider two surfaces $S_{1}$ (plane) and $S_{2}$ (hemispherical) bounded by common closed path $l$. Let an instant of time ' $t$ ', the current is ' $I$ '.
Applying Ampere's law for surface $\mathrm{S}_{1}$, we get
$\oint \overrightarrow{\boldsymbol{B}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{l}}=\mu_{0} I \ldots \ldots . .(1)$
Ampere's law for surface $S_{2}$ is given by,


$$
\begin{equation*}
\oint \vec{B} \cdot \boldsymbol{d} \vec{l}=0 . \tag{2}
\end{equation*}
$$

Since, the dielectric current inside the capacitor is zero. We see that equation (1) \& (2) are contradict to each other which is impossible. So, to remove this controversy, Maxwell introduced the modification in Ampere's law.


What is displacement current? [2015-16, 2016-17, 2018-19, 2020-21]
Or
Explain the concept of displacement current and show how it leads to the modification of Ampere's law.[2016-17, 2018-19, 2020-21]

## Displacement current

The concept of displacement current was introduced to resolve the paradox of charging capacitor. Maxwell proposed that


It means that changing electric field is equivalent to a current which flows as long as the electric field is changing. This equivalent current in vacuum or dielectric produces same magnetic field as the conduction current in conductor and is known as Displacement current .

Therefore, Maxwell introduced a factor $\varepsilon_{0} \frac{d \phi_{E}}{d t}$ to add in equation (1) instead of ' $I$ ' which is denoted as ' $I_{d}$ ' (displacement current). So,

$$
\oint \overrightarrow{\boldsymbol{B}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{l}}=\mu_{0}\left(I+\varepsilon_{0} \frac{d \phi_{E}}{d t}\right)
$$

## Modification of Ampere's law

Ampere's law in differential form is given by

$$
\begin{equation*}
\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J} \tag{1}
\end{equation*}
$$

The above equation stands only for steady state current but for time varying fields the current density should be modified. Thus taking divergence of both sides of above equation we get

$$
\vec{\nabla} \cdot(\vec{\nabla} \times \vec{H})=\vec{\nabla} \cdot \vec{J}
$$

$$
\begin{align*}
& \text { Since } \overrightarrow{\boldsymbol{\nabla}} \cdot(\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}})=0 \quad \text { (i.e. Divergence of curl (DC) is zero) } \\
& \text { Therefore, } \quad \overrightarrow{\boldsymbol{\nabla}} \cdot \vec{J}=0 \ldots \ldots . . . . . . . \text { (2) } \tag{2}
\end{align*}
$$

But according to the equation of continuity,
$\vec{\nabla} \cdot \vec{J}=-\frac{\partial \rho}{\partial t}$.

Maxwell realized the situation and suggested that the definition of total current density is incomplete and suggested to add something to J, such that equation 1 becomes
$\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\overrightarrow{\boldsymbol{J}}+\overrightarrow{\boldsymbol{J}}^{\prime}$
Now taking divergence of both the sides we get

$$
\begin{align*}
\vec{\nabla} \cdot(\vec{\nabla} \times \vec{H}) & =\nabla \cdot\left(\vec{J}+\vec{J}^{\prime}\right)  \tag{4}\\
\vec{\nabla} \cdot\left(\vec{J}+\vec{J}^{\prime}\right) & =0
\end{align*}
$$

$$
\vec{\nabla} . \vec{J}=-\vec{\nabla} . \vec{j}^{\prime}
$$

Putting the value of $\vec{\nabla}$. $\overrightarrow{\boldsymbol{J}}$ from equation (3) we get

## Or

$$
\vec{\nabla} \cdot \vec{J}^{\prime}=\frac{\partial \rho}{\partial t}
$$

From Maxwell's first equation,

$$
\overrightarrow{\boldsymbol{\nabla}} . \boldsymbol{D}=\rho
$$

Therefore

$$
\begin{aligned}
& \overrightarrow{\boldsymbol{\nabla}} \cdot \vec{J}^{\prime}=\frac{\partial}{\partial t}(\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}) \\
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{J}}^{\prime}=\overrightarrow{\boldsymbol{\nabla}} \cdot\left(\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}\right)
\end{aligned}
$$

Hence

$$
\vec{J}^{\prime}=\frac{\partial \vec{D}}{\partial t}
$$

Substituting the value of $J^{\prime}$ in equation 4 , we get
$\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t} \quad\{$ Modified Ampere's Law $\}$
where $\frac{\partial \vec{D}}{\partial t}$ is displacement current density which arises when the electric displacement vector $\overrightarrow{\boldsymbol{D}}$ is changing with time. It flows until the electric field is changing with time. In this way displacement current leads to the modification of Ampere's law.

What is the difference between conduction current and displacement current?
[2017-18]

## Difference between conduction and displacement current

| S. No. | Conduction current | Displacement current |
| :--- | :--- | :--- |
| 1. | The electric current carried <br> by conductors due to flow of <br> charge is called conduction <br> current | The electric current due to <br> changing electric field is called <br> displacement current |
| 2. | It exist even if flow of <br> electron is at uniform rate. | state. |
| 3. | $I_{c}=\mathrm{V} / \mathrm{R}$ |  |$\quad I_{d}=\varepsilon_{0} \frac{d \phi_{E}}{d t}$|  |
| :--- |

## Lecture 12 <br> ELECTROMAGNETIC FIELD THEORY



## CONTENTS

> Maxwell's equations $\longrightarrow$ Differential form
> Physical Significance of Maxwell's equations

Derive Maxwell's equation in differential form. Give physical significance of each equation. [2017-18]
or
Write Maxwell's equation in integral and differential form and explain their physical significance with their proofs. [2018-19]

## Maxwell's equations

## Maxwell's equation in Differential form:

## Maxwell's first equation

According to Gauss law in electrostatics
$\oint \overrightarrow{\boldsymbol{E}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=\frac{q}{\varepsilon_{0}}=\frac{1}{\varepsilon_{0}} \int \rho d V$

$$
\oint \varepsilon_{0} \overrightarrow{\boldsymbol{E}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=\int \rho d V
$$

$$
\oint \overrightarrow{\boldsymbol{D}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=\int \rho d V \quad\left(\text { Because } \overrightarrow{\boldsymbol{D}}=\varepsilon_{0} \overrightarrow{\boldsymbol{E}}\right)
$$

where $\overrightarrow{\boldsymbol{D}}$ is called electric displacement vector.

Using Gauss divergence theorem, we get
$\oint \overrightarrow{\boldsymbol{D}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{D}}) d V=\int \rho d V$
Therefore,

$$
\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{D}}-\rho) d V=0
$$

Since the equation is true for any volume the integrand must vanish, thus

$$
\begin{gathered}
(\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}-\rho)=0 \\
\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}=\rho
\end{gathered}
$$

This is Maxwell's first equation.

## Maxwell's second equation

According to Gauss law in magnetostatics

$$
\oint \vec{B} \cdot d \overrightarrow{\boldsymbol{S}}=0
$$

Using Gauss divergence theorem, we get

$$
\oint \overrightarrow{\boldsymbol{B}} \cdot d \overrightarrow{\boldsymbol{S}}=\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{B}}) d V
$$

Therefore,

$$
\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{B}}) d V=0
$$

The integrand should vanish for the surface boundary as the volume is arbitrary
i.e.
$\vec{\nabla} \cdot \vec{B}=0$

## Maxwell's third equation

According to Faraday's law of electromagnetic induction

$$
e m f=-\frac{d \emptyset_{B}}{d t}
$$

By the definition of emf we know

$$
e m f=\oint \overrightarrow{\boldsymbol{E}} \cdot d \overrightarrow{\boldsymbol{l}}
$$

And by the definition of magnetic flux we know

$$
\emptyset_{B}=\int \vec{B} \cdot d \vec{S}
$$

where S is the surface bounded by the circuit.

Thus, we have

$$
\oint \overrightarrow{\boldsymbol{E}} \cdot d \overrightarrow{\boldsymbol{l}}=-\frac{d}{d t} \int \overrightarrow{\boldsymbol{B}} \cdot d \overrightarrow{\boldsymbol{S}}
$$

Since the surface $S$ does not change its shape or position with time, we can write above equation as

$$
\oint \vec{E} \cdot d \vec{l}=-\int \frac{\partial \vec{B}}{\partial t} \cdot d \vec{S}
$$

Using Stoke's theorem, we get
$\oint \vec{E} \cdot d \vec{l}=\int(\vec{\nabla} \times \vec{E}) \cdot d \vec{S}=-\int \frac{\partial \vec{B}}{\partial t} \cdot d \vec{S}$

$$
\int\left(\vec{\nabla} \times \overrightarrow{\boldsymbol{E}}+\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t}\right) \cdot d \overrightarrow{\boldsymbol{S}}=\mathbf{0}
$$

This equation must hold for any arbitrary surface, thus the integrand should vanish

$$
\vec{\nabla} \times \overrightarrow{\boldsymbol{E}}+\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t}=0
$$

Or

$$
\vec{\nabla} \times \overrightarrow{\boldsymbol{E}}=-\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t}
$$

This is Maxwell's third equation.

## Maxwell's fourth equation

According to Ampere's circuital law

$$
\oint \overrightarrow{\boldsymbol{H}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{l}}=I\left[\oint \overrightarrow{\boldsymbol{B}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{l}}=\mu_{0} I \quad \text { and } B=\mu_{0} H\right]
$$

We also know

$$
\begin{gathered}
I=\int \vec{J} \cdot d \vec{S} \\
\oint \vec{H} \cdot d \vec{l}=\int \vec{J} \cdot d \vec{S}
\end{gathered}
$$

So,

Using Stoke's theorem, we get

$$
\oint \vec{H} \cdot d \vec{l}=\int(\vec{\nabla} \times \vec{H}) \cdot d \vec{S}=\int \vec{J} \cdot d \vec{S}
$$

$$
\int(\vec{\nabla} \times \vec{H}-\vec{J}) \cdot d \vec{S}=0
$$

For any arbitrary surface, this equation must hold. Thus, the integrand should vanish

$$
\begin{equation*}
\vec{\nabla} \times \vec{H}-\vec{J}=0 \tag{1}
\end{equation*}
$$

$\vec{\nabla} \times \vec{H}=\overrightarrow{\boldsymbol{J}}$.
This equation derived on the basis of ampere's law stands only for steady state current but for time varying fields the current density should be modified. Thus taking divergence of both sides of above equation we get,

$$
\vec{\nabla} \cdot(\vec{\nabla} \times \vec{H})=\vec{\nabla} \cdot \vec{J}
$$

$$
\begin{equation*}
\vec{\nabla} \cdot \vec{J}=0 \ldots \ldots . . . . . . . . . . .(2) \quad[\text { since }, \vec{\nabla} \cdot(\vec{\nabla} \times \overrightarrow{\boldsymbol{H}})=0] \tag{2}
\end{equation*}
$$

But by equation of continuity, we have
$\vec{\nabla} . \vec{J}=-\frac{\partial \rho}{\partial t}$.

Maxwell realized the situation and suggested that the definition of total current density is incomplete and suggested to add something to J , such that equation 1 becomes

$$
\begin{equation*}
\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J}+\vec{J}^{\prime} . \tag{4}
\end{equation*}
$$

Now taking divergence of both the sides we get

$$
\vec{\nabla} \cdot(\vec{\nabla} \times \vec{H})=\nabla \cdot\left(\vec{J}+\vec{J}^{\prime}\right)
$$

Or

$$
\begin{aligned}
& \vec{\nabla} \cdot\left(\vec{J}+\vec{J}^{\prime}\right)=0 \\
& \quad \vec{\nabla} \cdot \vec{J}=-\vec{\nabla} \cdot \vec{J}^{\prime}
\end{aligned}
$$

Putting the value of $\vec{\nabla}$. $\vec{J}$ from equation (3) we get
Or

$$
\vec{\nabla} \cdot \vec{J}^{\prime}=\frac{\partial \rho}{\partial t}
$$

But we know from Maxwell's first equation,

$$
\begin{gathered}
\vec{\nabla} \cdot \vec{D}=\rho \\
\vec{\nabla} \cdot \vec{J}^{\prime}=\frac{\partial}{\partial t}(\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}) \\
\vec{\nabla} \cdot \vec{J}^{\prime}=\vec{\nabla} \cdot\left(\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}\right)
\end{gathered}
$$

Hence

$$
\vec{J}^{\prime}=\frac{\partial \overrightarrow{\mathbf{D}}}{\partial t}
$$

Substituting the value of $J^{\prime}$ in equation 4 , we get

$$
\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}
$$

This is Maxwell's fourth equation. The second term arises when the electric displacement vector $\overrightarrow{\boldsymbol{D}}$ is changing with time and is called displacement current density, $\overrightarrow{\boldsymbol{J}}$ is called conduction current density.

## Maxwell's equations

## Maxwell's equation in Integral form

## (1) First Equation: -

We know Maxwell's first equation in differential form

$$
\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}=\rho
$$

Integrating over an entire volume we get

$$
\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{D}}) d V=\int \rho d V
$$

Applying Gauss divergence theorem on LHS we get

$$
\oint \vec{D} \cdot d \overrightarrow{\boldsymbol{S}}=\int \rho d V
$$

$$
\begin{gathered}
\oint\left(\varepsilon_{0} \vec{E}\right) \cdot d \vec{S}=\int \rho d V \\
\varepsilon_{0} \oint \vec{E} \cdot d \vec{S}=\int \rho d V \\
\oint \vec{E} \cdot d \vec{S}=\frac{1}{\varepsilon_{0}} \int \rho d V=\frac{q}{\varepsilon_{0}} \\
\oint \vec{E} \cdot d \vec{S}=\frac{q}{\varepsilon_{0}}
\end{gathered}
$$

## (2) Second Equation:

We know Maxwell's second equation in differential form

$$
\vec{\nabla} \cdot \vec{B}=0
$$

Integrating over an entire volume we get

$$
\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{B}}) d V=0
$$

Applying Gauss divergence theorem on LHS we get

$$
\oint \vec{B} \cdot d \vec{S}=0
$$

## (3) Third Equation:

We know Maxwell's third equation in differential form

$$
\vec{\nabla} \times \overrightarrow{\boldsymbol{E}}=-\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t}
$$

Integrating over an open surface area we get

$$
\int(\vec{\nabla} \times \vec{E}) \cdot d \vec{S}=-\int \frac{\partial \vec{B}}{\partial t} \cdot d \vec{S}
$$

Applying Stoke's theorem on LHS, $\int(\vec{\nabla} \times \vec{E}) \cdot d \vec{S}=\oint \vec{E} \cdot d \overrightarrow{\boldsymbol{l}}$ we get,

$$
\begin{aligned}
\oint \overrightarrow{\boldsymbol{E}} \cdot d \vec{l} & =-\int \frac{\partial \vec{B}}{\partial t} \cdot d \overrightarrow{\boldsymbol{S}} \\
\oint \overrightarrow{\boldsymbol{E}} \cdot d \vec{l} & =-\frac{d}{d t} \int \overrightarrow{\boldsymbol{B}} \cdot d \overrightarrow{\boldsymbol{S}}
\end{aligned}
$$

## (4) Fourth Equation:

We know

$$
\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}
$$

Integrating over an open surface area we get

$$
\begin{gathered}
\int(\vec{\nabla} \times \vec{H}) \cdot d \vec{S}=\int\left(\vec{J}+\frac{\partial \overrightarrow{\mathrm{D}}}{\partial t}\right) \cdot d \vec{S} \\
\int(\vec{\nabla} \times \vec{H}) \cdot d \vec{S}=\int \vec{J} \cdot d \vec{S}+\int \frac{\partial \overrightarrow{\mathrm{D}}}{\partial t} \cdot d \vec{S}
\end{gathered}
$$

$$
\int(\vec{\nabla} \times \vec{H}) \cdot d \vec{S}=\int \vec{J} \cdot d \vec{S}+\int \frac{\partial \vec{D}}{\partial t} \cdot d \vec{S}
$$

Applying Stoke's theorem on LHS we get

$$
\oint \overrightarrow{\boldsymbol{H}} \cdot d \vec{l}=I+\int \frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t} \cdot d \vec{S}
$$

It is modified Ampere`s law. It states that any current carrying conductor as well as time varying electric field produces a magnetic field.

## Physical Significance of Maxwell's equation

## First Equation:

It is Gauss law in electrostatics. It states that the surface integral of electric field over any closed surface area is equal to $\frac{1}{\varepsilon_{0}}$ times of net charge enclosed by that surface.

## Second Equation:

## Ques: Show that magnetic monopoles does not exist?[2020-21]

It is Gauss law in magneto statics. It states that there is no existence of magnetic monopoles or the net magnetic flux through any closed surface area is zero. It also signifies that magnetic field lines are closed curves.

## Physical Significance of Maxwell's equation

## Third Equation:

It is Faraday`s law in electromagnetic induction. It states that induced emf around any closed path is equal to the negative rate of change of magnetic flux bounded by the surface w.r.t. time. i.e. any changing magnetic field produces a electric field.

## Fourth Equation:

It is modified Ampere`s law. It states that any current carrying conductor as well as time varying electric field produces a magnetic field.

## Derive Coulomb`s law from Maxwell's first equation. [2018-19]

Maxwell's first equation, $\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{D}}=\rho$ Integrating over an entire volume we get

$$
\int(\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{D}}) d V=\int \rho d V
$$

Applying Gauss divergence theorem on LHS we get

$$
\begin{gathered}
\oint \overrightarrow{\boldsymbol{D}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=\int \rho d V \\
\oint\left(\varepsilon_{0} \overrightarrow{\boldsymbol{E}}\right) \cdot d \overrightarrow{\boldsymbol{S}}=\int \rho d V
\end{gathered}
$$



$$
\oint \mathrm{E} d S=\frac{q}{\varepsilon_{0}} \quad \begin{gathered}
\varepsilon_{0} \oint \overrightarrow{\boldsymbol{E}} \cdot \boldsymbol{d} \overrightarrow{\boldsymbol{S}}=q \\
(\text { ' } \mathrm{q} \text { ' is point charge })
\end{gathered}
$$

For a positive point charge electric field will be radial outward and so it will along $d \vec{S}$

$$
\begin{aligned}
& \mathrm{E} \oint \boldsymbol{d} \boldsymbol{S}=\frac{q}{\varepsilon_{0}} \\
& \text { E. } 4 \pi r^{2}=\frac{q}{\varepsilon_{0}} \\
& \mathrm{E}=\frac{q}{4 \pi \varepsilon_{0} r^{2}}
\end{aligned}
$$

If test charge $q_{0}$ is placed at point P , then electrostatic force experienced by $q_{0}$ is given by

$$
\begin{gathered}
F=q_{0} E \\
F=\frac{q_{0} q}{4 \pi \varepsilon_{0} r^{2}}
\end{gathered}
$$

This is Coulomb`s law

## Derive equation of continuity from Maxwell`s fourth equation.

Maxwells fourth equation is, $\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}}=\overrightarrow{\boldsymbol{J}}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}$
Taking divergence on both sides we get

$$
\vec{\nabla} \cdot(\vec{\nabla} \times \overrightarrow{\boldsymbol{H}})=\vec{\nabla} \cdot\left(\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}\right)
$$

$\vec{\nabla} \cdot\left(\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}\right)=0 ; \quad[\operatorname{since} \overrightarrow{\boldsymbol{\nabla}} \cdot(\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}})=0]$

$$
\vec{\nabla} \cdot \vec{J}+\vec{\nabla} \cdot\left(\frac{\partial \vec{D}}{\partial t}\right)=0
$$

$$
\vec{\nabla} \cdot \vec{J}+\frac{\partial(\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}})}{\partial t}=0
$$

From Maxwell`s first equation we know,

$$
\vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}=\rho
$$

Putting the value of $\overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{D}}$,
we get, $\vec{\nabla} \cdot \vec{J}+\frac{\partial \rho}{\partial t}=0$

This is the equation of continuity.

## Lecture 13 ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Poynting theorem
> Poynting vector

State and deduce Poynting theorem for the flow of energy in an electromagnetic field. Also, discuss its physical significance.
[2016-17, 2019-20, 2020-21]
Or
Discuss the work-energy theorem for the flow of energy in an electromagnetic field.
[2018-19]

## Poynting theorem

Poynting theorem describes the flow of energy or power in an electromagnetic field during the propagation of uniform plane wave. Maxwell's third and fourth equations are given

$$
\begin{align*}
& \vec{\nabla} \times \overrightarrow{\boldsymbol{E}}=-\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t}  \tag{1}\\
& \overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}}=\overrightarrow{\boldsymbol{J}}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t} \tag{2}
\end{align*}
$$

Taking dot product of equation (1) with $\overrightarrow{\boldsymbol{H}}$ and equation (2) with $\overrightarrow{\boldsymbol{E}}$, we get

$$
\begin{gather*}
\overrightarrow{\boldsymbol{H}} \cdot(\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}})=-\overrightarrow{\boldsymbol{H}} \cdot \frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial \mathrm{t}}  \tag{3}\\
\overrightarrow{\boldsymbol{E}} \cdot(\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}})=\overrightarrow{\boldsymbol{E}} \cdot\left(\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}\right) \tag{4}
\end{gather*}
$$

Subtracting equation (4) from equation (3) we get

$$
\begin{equation*}
\overrightarrow{\boldsymbol{H}} \cdot(\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}})-\overrightarrow{\boldsymbol{E}} \cdot(\vec{\nabla} \times \overrightarrow{\boldsymbol{H}})=-\overrightarrow{\boldsymbol{H}} \cdot \frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial \mathrm{t}}-\overrightarrow{\boldsymbol{E}} \cdot \overrightarrow{\boldsymbol{J}}-\overrightarrow{\boldsymbol{E}} \cdot \frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t} \tag{5}
\end{equation*}
$$

Using vector identity, $\overrightarrow{\boldsymbol{\nabla}} \cdot(\overrightarrow{\boldsymbol{E}} \times \overrightarrow{\boldsymbol{H}})=\overrightarrow{\boldsymbol{H}} \cdot(\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}})-\overrightarrow{\boldsymbol{E}} \cdot(\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}) \ldots$ (6)

From equations (5) and (6), we get

$$
\begin{equation*}
\vec{\nabla} \cdot(\overrightarrow{\boldsymbol{E}} \times \overrightarrow{\boldsymbol{H}})=-\overrightarrow{\boldsymbol{H}} \cdot \frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial \mathrm{t}}-\overrightarrow{\boldsymbol{E}} \cdot \vec{J}-\overrightarrow{\boldsymbol{E}} \cdot \frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t} \tag{7}
\end{equation*}
$$

on rearranging, $\overrightarrow{\boldsymbol{\nabla}} \cdot(\overrightarrow{\boldsymbol{E}} \times \overrightarrow{\boldsymbol{H}})=-\overrightarrow{\boldsymbol{E}} \cdot \overrightarrow{\boldsymbol{J}}-\overrightarrow{\boldsymbol{E}} \cdot \frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t}-\overrightarrow{\boldsymbol{H}} \cdot \frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial \mathrm{t}}$

$$
\begin{equation*}
\text { or } \quad-\vec{\nabla} \cdot(\overrightarrow{\boldsymbol{E}} \times \overrightarrow{\boldsymbol{H}})=\overrightarrow{\boldsymbol{E}} \cdot \vec{J}+\overrightarrow{\boldsymbol{E}} \cdot \frac{\partial \vec{D}}{\partial t}+\overrightarrow{\boldsymbol{H}} \cdot \frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial \mathrm{t}} \tag{8}
\end{equation*}
$$

Now, since $\overrightarrow{\boldsymbol{D}}=\varepsilon \overrightarrow{\boldsymbol{E}}$ and $\overrightarrow{\boldsymbol{B}}=\mu \overrightarrow{\boldsymbol{H}}$

$$
\vec{E} \cdot \frac{\partial \vec{D}}{\partial t}=\varepsilon \vec{E} \cdot \frac{\partial \vec{E}}{\partial t}=\frac{1}{2} \varepsilon \frac{\partial(\vec{E} \cdot \vec{E})}{\partial t}=\frac{\partial}{\partial t}\left[\frac{1}{2} \varepsilon\left(E^{2}\right)\right]
$$

## Similarly,

$$
\vec{H} \cdot \frac{\partial \vec{B}}{\partial t}=\mu \vec{H} \cdot \frac{\partial \vec{H}}{\partial t}=\frac{1}{2} \mu \frac{\partial(\overrightarrow{\boldsymbol{H}} \cdot \overrightarrow{\boldsymbol{H}})}{\partial t}=\frac{\partial}{\partial t}\left[\frac{1}{2} \mu\left(H^{2}\right)\right]
$$

Making these substitutions in equation (8), we get

$$
-\vec{\nabla} \cdot(\vec{E} \times \vec{H})=\vec{E} \cdot \vec{J}+\frac{\partial}{\partial t}\left[\frac{1}{2} \varepsilon\left(E^{2}\right)\right]+\frac{\partial}{\partial t}\left[\frac{1}{2} \mu\left(H^{2}\right)\right]
$$

Taking the volume integral of the above equation we get
$-\int \vec{\nabla} \cdot(\vec{E} \times \vec{H}) d V=\int(\vec{E} \cdot \vec{J}) d V+\int \frac{\partial}{\partial t}\left[\frac{1}{2} \varepsilon\left(E^{2}\right)\right] d V+\int \frac{\partial}{\partial t}\left[\frac{1}{2} \mu\left(H^{2}\right)\right] d V$

Applying Gauss divergence theorem on LHS, we get
$-\int(\vec{E} \times \vec{H}) \cdot \overrightarrow{d S}=\int(\vec{E} \cdot \vec{J}) d V+\frac{d}{d t} \int \frac{1}{2}\left[\varepsilon\left(E^{2}\right)+\mu\left(H^{2}\right)\right] d V$
Or $\int(\vec{E} \cdot \overrightarrow{J)} d V=-\frac{d}{d t} \underbrace{\int \frac{1}{2}\left[\varepsilon\left(E^{2}\right)+\mu\left(H^{2}\right)\right] d V-\int(\vec{E} \times \vec{H})}_{\text {I term }} \cdot \overrightarrow{d S}$
Above equation represent work-energy theorem and is called Poynting theorem. In above equation,
(i) The first term represents the total power dissipated in a volume ' $V$ '.
(ii) The second term represents the sum of energy stored in electric field $\left\{\frac{1}{2} \int\left[\varepsilon\left(E^{2}\right)\right\}\right.$ and in magnetic field $\left.\left\{\frac{1}{2} \int \mu\left(H^{2}\right)\right] d V\right\}$.
(iii) The third term represents the rate at which the energy is carried out of volume ' V ', across its boundary surface by electromagnetic wave.
> Physical significance : Above (i), (ii), and (iii) statement represent the physical significance of Poynting theorem or we can say that physical significance of Poynting theorem is 'conservation of energy', as it is clear from equation (9) that the work done on the charge by an EM force is equal to decrease in energy stored in field, less than the energy which flowed out through surface.

What is Poynting vector? [2016-17, 2018-19]

## Poynting vector

The amount of field energy passing through unit area of the surface per unit time is called Poynting vector. It is represented by

$$
\vec{S}=\vec{E} \times \vec{H}
$$

Since, $\overrightarrow{\boldsymbol{E}}$ and $\overrightarrow{\boldsymbol{H}}$ are perpendicular to each other, thus the magnitude of $\overrightarrow{\boldsymbol{S}}$ is given by
$\overrightarrow{\boldsymbol{S}}=\boldsymbol{E} \boldsymbol{H} \quad[\overrightarrow{\boldsymbol{S}}=\overrightarrow{\boldsymbol{E}} \times \overrightarrow{\boldsymbol{H}}=\mathrm{EH} \sin \theta]$
The direction of $\overrightarrow{\boldsymbol{S}}$ is along the direction of wave propagation.

## What is the unit and dimension of Poynting vector?

Unit of Poynting vector is Joule/ ( $\mathrm{m}^{2}$.sec) or Watt/m $\mathrm{m}^{2}$.
Dimensions of Poynting vector are

$$
\begin{gathered}
\frac{\text { Energy }}{\text { Area. Time }}=\frac{M L^{2} T^{-2}}{L^{2} T} \\
=M T^{-3}
\end{gathered}
$$

## Lecture 14 <br> ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Electromagnetic wave in free space
$>$ Electromagnetic wave in non conducting medium
$>$ Transverse nature of electromagnetic waves

Deduce four Maxwell's equation in free space.[2019-20, 2020-21]

## Electromagnetic wave in free space

Maxwell's equations are

$$
\begin{align*}
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{D}}=\rho  \tag{1}\\
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{B}}=0 \tag{2}
\end{align*}
$$

$$
\begin{equation*}
\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}}=-\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t} \tag{3}
\end{equation*}
$$

$$
\begin{equation*}
\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J}+\frac{\partial \overrightarrow{\boldsymbol{D}}}{\partial t} . \tag{4}
\end{equation*}
$$

Vacuum means that there is no free charge in the region so there is no current produced, i.e., no conduction current. So, we have
$\rho=0, \overrightarrow{\boldsymbol{J}}=0, \overrightarrow{\boldsymbol{D}}=\varepsilon_{0} \overrightarrow{\boldsymbol{E}}, \overrightarrow{\boldsymbol{B}}=\mu_{0} \overrightarrow{\boldsymbol{H}}$
Making these substitutions in above equations. Then, Maxwell's equations in free space is given as

$$
\begin{align*}
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{E}}=0  \tag{5}\\
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{H}}=0  \tag{6}\\
& \vec{\nabla} \times \overrightarrow{\boldsymbol{E}}=-\mu_{0} \frac{\partial \overrightarrow{\boldsymbol{H}}}{\partial t}  \tag{7}\\
& \vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\varepsilon_{0} \frac{\partial \overrightarrow{\boldsymbol{E}}}{\partial t} \tag{8}
\end{align*}
$$

Derive the electromagnetic wave equations in free space. Prove that the electromagneticwavespropagatewithspeedof light in free space.
[2017-18]

## Or

Derive the equation for the propagation of plane electromagnetic wave in free space. Show that the velocity of plane electromagnetic wave in free space is given by $c=1 / \sqrt{\mu_{0} \epsilon_{0}}$.
[2015-16, 2018-19]

Maxwell`s equations in free space is given as

$$
\begin{align*}
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{E}}=0 \\
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{H}}=0 \\
& \overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}}=-\mu_{0} \frac{\partial \overrightarrow{\boldsymbol{H}}}{\partial t}  \tag{3}\\
& \vec{\nabla} \times \vec{H}=\varepsilon_{0} \frac{\partial \vec{E}}{\partial t} \tag{4}
\end{align*}
$$

Taking curl of equation (3), we get

$$
\begin{aligned}
& \vec{\nabla} \times\left(\vec{\nabla} \times \overrightarrow{E)}=-\vec{\nabla} \times\left(\mu_{0} \frac{\partial \vec{H}}{\partial t}\right)\right. \\
& \Rightarrow \vec{\nabla} \times\left(\vec{\nabla} \times \overrightarrow{E)}=-\mu_{0} \frac{\partial}{\partial t}(\vec{\nabla} \times \vec{H}) \cdots \ldots(5)\right.
\end{aligned}
$$

Putting the value of $\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}}$ in equation (5) from equation (4) we get

$$
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\mu_{0} \frac{\partial}{\partial t}\left(\varepsilon_{0} \frac{\partial \vec{E}}{\partial t}\right)
$$

$$
\begin{equation*}
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\mu_{0} \varepsilon_{0} \frac{\partial^{2} E}{\partial t^{2}} . \tag{6}
\end{equation*}
$$

Using vector identity (vector triple product),

$$
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=\vec{\nabla}(\vec{\nabla} \cdot \vec{E})-\nabla^{2} \vec{E}
$$

From equation (1) we have $\vec{\nabla} \cdot \vec{E}=0$,
Therefore

$$
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\nabla^{2} \vec{E}
$$

Putting this in equation (6), we get

$$
-\nabla^{2} \vec{E}=-\mu_{0} \varepsilon_{0} \frac{\partial^{2} E}{\partial t^{2}}
$$

$$
\nabla^{2} \vec{E}=\mu_{0} \varepsilon_{0} \frac{\partial^{2} E}{\partial t^{2}}
$$

Similarly, $\nabla^{2} \vec{H}=\mu_{0} \varepsilon_{0} \frac{\partial^{2} H}{\partial t^{2}} \ldots . .$. (8)
By general wave equation we have

$$
\begin{equation*}
\nabla^{2} \psi=\frac{1}{v^{2}} \frac{\partial^{2} \psi}{\partial t^{2}} \tag{9}
\end{equation*}
$$

where $\psi$ is the wave function which propagates with velocity $v$. Thus, comparing equations (7) \& (8) with the above equation (9), we observe that field vectors E \& H propagate as waves in free space and the velocity of propagation is

$$
\begin{aligned}
& v=\frac{1}{\sqrt{\mu_{0} \varepsilon_{0}}} \\
& =\frac{1}{\sqrt{4 \pi \times 10^{-7} \times 8.85 \times 10^{-12}}} \\
& =2.99 \times 10^{8} \mathrm{~m} / \mathrm{s}=c(\text { Speed of light })
\end{aligned}
$$

Thus, EM waves travel in free space with speed of light.So, equations (7) and (8) may be written as

$$
\nabla^{2} \overrightarrow{\boldsymbol{E}}-\frac{1}{c^{2}} \frac{\partial^{2} \overrightarrow{\boldsymbol{E}}}{\partial t^{2}}=0 \& \nabla^{2} \overrightarrow{\boldsymbol{H}}-\frac{1}{c^{2}} \frac{\partial^{2} \overrightarrow{\boldsymbol{H}}}{\partial t^{2}}=0
$$

Write the Maxwell's equation in non conducting medium. Also, derive electromagnetic wave equation in non conducting medium.

## Electromagnetic wave in non conducting medium

Maxwell's equations are

$$
\begin{align*}
\vec{\nabla} \cdot \vec{D} & =\rho  \tag{1}\\
\vec{\nabla} \cdot \vec{B} & =0 \tag{2}
\end{align*}
$$

$\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}}=-\frac{\partial \overrightarrow{\boldsymbol{B}}}{\partial t}$
$\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}=\vec{J}+\frac{\partial \vec{D}}{\partial t}$

Non conducting medium means non-charged, current free dielectric are non-conducting media. So, we have

$$
\rho=0, \vec{J}=0, \vec{D}=\varepsilon \vec{E}, \vec{B}=\mu \vec{H}
$$

Making these substitutions in above equations, we get Maxwell`s equations in non conducting medium

$$
\begin{align*}
& \overrightarrow{\boldsymbol{\nabla}} \cdot \overrightarrow{\boldsymbol{E}}=0  \tag{5}\\
& \vec{\nabla} \cdot \overrightarrow{\boldsymbol{H}}=0  \tag{6}\\
& \overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}}=-\mu \frac{\partial \overrightarrow{\boldsymbol{H}}}{\partial t}  \tag{7}\\
& \overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{H}}=\varepsilon \frac{\partial \overrightarrow{\boldsymbol{E}}}{\partial t} \tag{8}
\end{align*}
$$

Taking curl of equation (7), we get

$$
\begin{array}{r}
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=\vec{\nabla} \times\left(\mu \frac{\partial \vec{H}}{\partial t}\right) \\
\Rightarrow \vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\mu \frac{\partial}{\partial t}(\vec{\nabla} \times \vec{H}) \cdots \ldots .(9) \tag{9}
\end{array}
$$

Putting the value of $\vec{\nabla} \times \overrightarrow{\boldsymbol{H}}$ in equation (9) from equation (8) we get

$$
\begin{gather*}
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\mu \frac{\partial}{\partial t}\left(\varepsilon \frac{\partial \vec{E}}{\partial t}\right) \\
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\mu \varepsilon \frac{\partial^{2} E}{\partial t^{2}} \ldots \ldots . .(1 \tag{10}
\end{gather*}
$$



$$
\vec{\nabla} \times(\vec{\nabla} \times \vec{E})=\vec{\nabla}(\vec{\nabla} \cdot \vec{E})-\nabla^{2} \vec{E}
$$

From equation (1) we have $\vec{\nabla} \cdot \vec{E}=0$,
Therefore $\quad \vec{\nabla} \times(\vec{\nabla} \times \vec{E})=-\nabla^{2} \vec{E}$
Putting this in equation (6), we get

$$
\begin{aligned}
& \nabla^{2} \vec{E}=\mu \varepsilon \frac{\partial^{2} E}{\partial t^{2}} \\
& \nabla^{2} \vec{H}=\mu \varepsilon \frac{\partial^{2} H}{\partial t^{2}}
\end{aligned}
$$

Similarly,

Prove that electromagnetic waves are transverse in nature.

> [2017-18, 2018-19]
or
Show that electric and magnetic vectors are normal to the direction of propagation of electromagnetic wave. [2020-21]
or
Show that vector $E, H$ and direction of propagation form a set of orthogonal vectors.
[2016-17]

## Transverse nature of electromagnetic waves

The equation for electric and magnetic field in free space are given by

$$
\nabla^{2} \vec{E}=\mu_{0} \varepsilon_{0} \frac{\partial^{2} E}{\partial t^{2}} ; \nabla^{2} \vec{H}=\mu_{0} \varepsilon_{0} \frac{\partial^{2} H}{\partial t^{2}}
$$

The solution of above equations may be written as:

$$
\begin{aligned}
\overrightarrow{\boldsymbol{E}}(r, t) & =\boldsymbol{E}_{0} e^{i(\overrightarrow{\boldsymbol{k}} \cdot \vec{r}-\omega t)} \\
\overrightarrow{\boldsymbol{H}}(r, t) & =\boldsymbol{H}_{0} e^{i(\overrightarrow{\boldsymbol{k}} \cdot \vec{r}-\omega t)}
\end{aligned}
$$

where $E_{0}$ and $H_{0}$ are amplitudes of electric \& magnetic fields and k ispropagation vector defined as

$$
\overrightarrow{\boldsymbol{k}}=k \hat{n}=\frac{2 \pi}{\lambda} \hat{n}=\frac{2 \pi v}{c} \hat{n}=\frac{\omega}{c} \hat{n}
$$

Here $\hat{n}$ is a unit vector in the direction of wave propagation.
Since there are no angular coordinates in $\overrightarrow{\boldsymbol{E}}$ and $\overrightarrow{\boldsymbol{H}}$, therefore

$$
\begin{gather*}
\vec{\nabla}=\frac{\partial}{\partial r} \\
\frac{\partial \overrightarrow{\boldsymbol{E}}}{\partial r}=\frac{\partial}{\partial r} \boldsymbol{E}_{\mathbf{0}} e^{i(\overrightarrow{\mathbf{k}} \cdot \overrightarrow{\boldsymbol{r}}-\omega t)}\left(\text { since } \frac{\partial\left(\boldsymbol{e}^{i x}\right)}{\partial x}=\boldsymbol{i} \boldsymbol{e}^{i x}\right) \\
=i k \boldsymbol{E}_{\mathbf{0}} e^{i(\overrightarrow{\boldsymbol{k}} \cdot \vec{r}-\omega t)} \\
=i k \overrightarrow{\boldsymbol{E}} \\
\Rightarrow \vec{\nabla}=\frac{\partial}{\partial r}=i k \ldots \ldots \ldots \ldots \ldots \ldots(1) \tag{1}
\end{gather*}
$$

Now

Now $\frac{\partial \overrightarrow{\boldsymbol{E}}}{\partial t}=\frac{\partial}{\partial t} \boldsymbol{E}_{\mathbf{0}} e^{i(\overrightarrow{\boldsymbol{k}} \cdot \overrightarrow{\boldsymbol{r}}-\omega t)}$

$$
\begin{gathered}
=(-i \omega) \boldsymbol{E}_{0} e^{i(\overrightarrow{\boldsymbol{k}} \cdot \overrightarrow{\boldsymbol{r}}-\omega t)} \\
=-i \omega \overrightarrow{\boldsymbol{E}} \\
\Rightarrow \quad \frac{\partial}{\partial t}=-i \omega \ldots \ldots \ldots \ldots(2)
\end{gathered}
$$

Making the substitutions for $\nabla$ in equation $\vec{\nabla} \cdot \overrightarrow{\boldsymbol{E}}=\mathbf{0}$, we get

$$
\begin{aligned}
i \overrightarrow{\boldsymbol{k}} \cdot \overrightarrow{\boldsymbol{E}} & =\mathbf{0} \\
\overrightarrow{\boldsymbol{k}} \cdot \overrightarrow{\boldsymbol{E}} & =0
\end{aligned}
$$

i.e.
$\overrightarrow{\boldsymbol{k}} \perp \overrightarrow{\boldsymbol{E}}$

Similarly making the substitutions for $\nabla$ in equation $\vec{\nabla} \cdot \vec{E}=0$ we get

$$
\begin{aligned}
i \overrightarrow{\boldsymbol{k}} \cdot \overrightarrow{\boldsymbol{H}} & =\mathbf{0} \\
\overrightarrow{\boldsymbol{k}} \cdot \overrightarrow{\boldsymbol{H}} & =0
\end{aligned}
$$

i.e.

$$
\begin{gathered}
\overrightarrow{\boldsymbol{k}} \perp \overrightarrow{\boldsymbol{H}} \\
\overrightarrow{\boldsymbol{k}} \perp \overrightarrow{\boldsymbol{H}} \perp \overrightarrow{\boldsymbol{E}}
\end{gathered}
$$

Now making the substitutions for $\nabla$ and $\frac{\partial}{\partial t}$ from eq (1) and (2) in equation
$\overrightarrow{\boldsymbol{\nabla}} \times \overrightarrow{\boldsymbol{E}}=-\mu_{0} \frac{\partial \overrightarrow{\boldsymbol{H}}}{\partial t}$, we get

$$
\begin{gathered}
i(\overrightarrow{\boldsymbol{k}} \times \overrightarrow{\boldsymbol{E}})=\mu_{0} i \omega \overrightarrow{\boldsymbol{H}} \\
k(\widehat{\boldsymbol{n}} \times \overrightarrow{\boldsymbol{E}})=\mu_{o} \omega \overrightarrow{\boldsymbol{H}} \\
\frac{k}{\mu_{\mathrm{o}} \omega}(\widehat{\boldsymbol{n}} \times \overrightarrow{\boldsymbol{E}})=\overrightarrow{\boldsymbol{H}} \\
\frac{1}{\mu_{\mathrm{o}} \mathrm{C}}(\widehat{\boldsymbol{n}} \times \overrightarrow{\boldsymbol{E}})=\overrightarrow{\boldsymbol{H}}\left[\text { since } \frac{\omega}{k}=c\right]
\end{gathered}
$$

From above equation, it is clear that field vector $\vec{H}$ is perpendicular to both $\vec{k}$ and $\vec{E}$. Thus, we see that electric field vectors E, magnetic field vectors H and the direction of propagation vector $\overrightarrow{\boldsymbol{k}}$ are mutually perpendicular to each other. This implies that electromagnetic waves are transverse in nature.

## Lecture 15 <br> ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Relation between electric and magnetic field
$>$ Energy and momentum of electromagnetic wave

What do you mean by impedance of a wave? [2019-20] or
What do you understand by characteristics impedance?

## Relation between electric and magnetic field

The ratio of magnitude of E to the magnitude of H is symbolized as $Z_{o}$ and has the dimensions of electric resistance.

We know

$$
\begin{gathered}
(\overrightarrow{\boldsymbol{k}} \times \overrightarrow{\boldsymbol{E}})=\mu_{0} \omega \overrightarrow{\boldsymbol{H}} \\
k E \sin 90^{\circ}=\mu_{0} \omega H \\
k E=\mu_{0} \omega H \\
E=\mu_{0} \frac{\omega}{k} H
\end{gathered}
$$

$$
\begin{array}{cc}
E=\mu_{0} c \boldsymbol{H} & \left(v=\frac{\omega}{k}=c\right) \\
E=\mu_{0} \frac{1}{\sqrt{\varepsilon_{0} \mu_{0}}} H & \left(c=\frac{1}{\sqrt{\mu_{0} \varepsilon_{0}}}\right) \\
E=\sqrt{\frac{\mu_{0}}{\varepsilon_{0}}} H &
\end{array}
$$

Therefore, $\quad Z_{o}=\left|\frac{\overrightarrow{\vec{E}}}{\overrightarrow{\vec{H}}}\right|=\frac{E_{o}}{H_{o}}=\sqrt{\frac{\mu_{o}}{\varepsilon_{o}}}$
 or wave impedance of free space.

$$
Z_{o}=376.72 \mathrm{ohms}
$$

Obtain the expression for total electromagnetic energy density.

## Energy density of plane EM wave in free space

The electric energy per unit volume or electric field energy density ' $u_{\mathrm{E}}$ ' is given by
$\mathrm{u}_{\mathrm{E}}=\frac{1}{2} \varepsilon_{0} E^{2}$
Similarly, the magnetic field energy density $U_{B}$ is given by

$$
\begin{equation*}
\mathrm{u}_{\mathrm{B}}=\frac{1}{2} \mu_{o} \mathrm{H}^{2} . \tag{2}
\end{equation*}
$$

both $\vec{E}$ and $\vec{B}$ is

$$
\begin{gather*}
\mathrm{u}=\mathrm{u}_{\mathrm{E}}+\mathrm{u}_{\mathrm{B}}=\frac{1}{2} \varepsilon_{0} E^{2}+\frac{1}{2} \mu_{o} \mathrm{H}^{2}, \\
\frac{E}{H}=\sqrt{\frac{\mu_{o}}{\varepsilon_{o}}} \quad \text { or } \quad \mathrm{H}=\mathrm{E} \sqrt{\frac{\varepsilon_{o}}{\mu_{o}}} \\
\mathrm{u}=\frac{1}{2} \varepsilon_{0} E^{2}+\frac{1}{2} \mu_{0} \frac{\varepsilon_{o}}{\boldsymbol{\mu}_{\boldsymbol{o}}} E^{2} \\
\mathrm{u}=\varepsilon_{0} E^{2} \ldots \ldots .(4) \\
\mathrm{u}=\mu_{o} \mathrm{H}^{2} \tag{5}
\end{gather*}
$$

Show that magnetic field and electric field of electromagnetic wave carries equal energies in free space.

We know that electric energy density, $\quad \mathrm{u}_{\mathrm{E}}=\frac{1}{2} \varepsilon_{0} E^{2}$
And magnetic energy density, $\mathrm{u}_{\mathrm{B}}=\frac{1}{2} \mu_{o} \mathrm{H}^{2}$.
Divide (1) from (2), $\quad \frac{\mathrm{u}_{\mathrm{E}}}{\mathrm{u}_{\mathrm{B}}}=\frac{\frac{1}{2} \varepsilon_{0} E^{2}}{\frac{1}{2} \mu_{0} H^{2}}=\frac{\varepsilon_{0} E^{2}}{\mu_{0} H^{2}}$

$$
\begin{array}{ll}
\frac{\mathrm{u}_{\mathrm{E}}}{\mathrm{u}_{\mathrm{B}}}=\frac{\varepsilon_{0} E^{2}}{\varepsilon_{0} E^{2}}=1 ; & {\left[\text { Since, } E=\sqrt{\frac{\mu_{0}}{\varepsilon_{0}}} H\right]} \\
\Rightarrow & \mathrm{u}_{\mathrm{E}}=\mathrm{u}_{\mathrm{B}}
\end{array}
$$

i.e., EM wave consists of equal electric and magnetic field energy.

Derive a relation between Poynting vector and total energy density.

Since,

$$
\vec{S}=\vec{E} \times \vec{H}
$$

If $\overrightarrow{\boldsymbol{E}}$ is along $\hat{x}$ - axis, $\overrightarrow{\boldsymbol{H}}$ is along $\hat{y}$ - axis. Then, $\overrightarrow{\boldsymbol{S}}$ is along $\hat{z}$ - axis

Therefore,

$$
\begin{aligned}
& \overrightarrow{\boldsymbol{S}}=E H(\hat{x} \times \hat{y})=E H \hat{z} \\
& \overrightarrow{\boldsymbol{S}}=\sqrt{\frac{\varepsilon_{0} x \varepsilon_{0}}{\mu_{0} \varepsilon_{0}}} E^{2} \hat{z}\left(\text { Since, } E=\sqrt{\frac{\mu_{0}}{\varepsilon_{0}}} H\right)
\end{aligned}
$$

$\overrightarrow{\boldsymbol{S}}=c \varepsilon_{0} E^{2} \hat{Z}$
Since, $c=\frac{1}{\sqrt{\mu_{0} \varepsilon_{0}}}$

$$
\vec{S}=\varepsilon_{0} E^{2} c \hat{z}=\mathrm{uc} \hat{z} ; \quad\left[\mathrm{u}=\varepsilon_{0} E^{2}\right]
$$

Thus, energy density associated with an electromagnetic wave in free space travels with speed equal to velocity of light with which the field vectors propagate.

## Lecture 16

## ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Radiation pressure (Resultant pressure)
> Skin depth
> Numericals

## What is radiation pressure?

## Radiation pressure (Resultant pressure)

> When electromagnetic radiation strikes to a surface area; its momentum gets changed. It means some momentum get transferred to the surface. If electromagnetic radiation is striking continuously to a surface, then the momentum of the surface will be continuously changed.
$>$ According to Newton`s second law the rate of change of momentum of the surface will be equal to the force exerted on the surface, i.e.

$$
\overrightarrow{\boldsymbol{F}}=\frac{d \overrightarrow{\boldsymbol{p}}}{d t}
$$

$>$ The force per unit area on the surface is equal to the pressure. Therefore, if electromagnetic radiation strikes to a surface, then it exerts a pressure on that surface, which is known as radiation pressure.

Derive a suitable expression for momentum and radiated pressure of an EM wave.

Or
Derive a suitable expression for energy density and radiated pressure of an EM wave.

Maxwell predicted that electromagnetic waves transport linear momentum in the direction of propagation. The momentum of the particle of mass $m$ moving with velocity $v$ is given by
$\overrightarrow{\boldsymbol{p}}=\boldsymbol{m} \overrightarrow{\boldsymbol{v}} \ldots \ldots \ldots \ldots$ (1)

According to Einstein's mass-energy relation,

$$
U=m c^{2} \quad \text { or } m=\frac{U}{c^{2}} \ldots \ldots \text { (2) }
$$

Using equation (2) in (1), then momentum density associated with EM wave is

$$
\overrightarrow{\boldsymbol{p}}=\frac{U}{c^{2}} \overrightarrow{\boldsymbol{v}}
$$

If electromagnetic wave propagating along $x$ direction, then $\vec{v}=c \hat{\imath}$
so,

$$
\vec{p}=\frac{U}{c} \hat{\imath} \Rightarrow p=\frac{U}{c}
$$

Let a plane EM wave incident normally on a perfectly absorbing surface having area A . Also let this surface area absorbs energy U in time t .

Then momentum transferred to the surface is
$p=\frac{U}{c}$

We know $S=\frac{U}{A t}$
where $S$ is magnitude of poynting vector.
Therefore,

$$
\begin{array}{ll}
p=\frac{S A t}{c} & \\
p=u A t & (S=u c)
\end{array}
$$

Where $u$ is energy density.
By Newton`s second law, $\quad F=\frac{d p}{d t}=u A$

Now, force exerted per unit area by EM wave is called radiation pressure, i.e., $P=\frac{F}{A}$
$P=\frac{1}{A} \frac{d p}{d t}$

Equation (5) gives required relation between radiation pressure and momentum. Now, Using equation (4) and (5)

$$
P=u
$$

i.e. Radiation pressure is equal to energy density.

What do you mean by depth of penetration? [2016-17, 2018-19]

## Skin depth (Depth of penetration)

When an EM wave propagate in a medium its amplitude decreases with the distance inside the medium from the surface. This phenomenon is known as attenuation. The amplitude of an EM wave at a depth x is given by

$$
E=E_{0} e^{-\alpha x}
$$

Where- $\mathrm{E}_{0}$ is amplitude of the wave at the surface of the medium, $\alpha$ is attenuation constant

The depth of penetration is defined as the depth in which the strength of electric field associated with the electromagnetic wave reduces to ${ }_{e}^{1}$ times to its initial value.


## What is the relation between skin depth and attenuation constant?

Since, $\quad E=E_{0} e^{-\alpha x} \quad \frac{E_{0}}{e}=E_{0} e^{-\alpha \delta}$

$$
\begin{aligned}
\text { At skin depth } x=\delta, \quad \begin{array}{c}
E \\
= \\
e^{-1}
\end{array}=e^{-\alpha \delta}
\end{aligned}
$$

$$
\delta=\frac{1}{\alpha}
$$

Thus,

$$
\text { skin depth }=\frac{1}{\text { attenuation constant }}
$$

Derive an expression of skin depth for good conductor and insulator. Also, show that skin depth for insulator does not depends on frequency of EM wave.

Attenuation constant ' $\alpha$ ' is given as

$$
\begin{equation*}
\alpha=\omega\left[\frac{\mu \varepsilon}{2}\left\{\left(1+\frac{\sigma^{2}}{\omega^{2} \varepsilon^{2}}\right)^{\frac{1}{2}}-1\right\}\right]^{\frac{1}{2}} \tag{1}
\end{equation*}
$$

For good conductors:

$$
\frac{\sigma}{\omega \varepsilon} \gg 1
$$

Therefore neglecting 1 with respect to $\frac{\sigma}{\omega \varepsilon}$ we get

$$
\begin{gathered}
\alpha=\omega\left\{\frac{\mu \varepsilon}{2}\left(\frac{\sigma}{\omega \varepsilon}-1\right)\right\}^{\frac{1}{2}} \\
\alpha=\omega\left(\frac{\mu \varepsilon}{2} \frac{\sigma}{\omega \varepsilon}\right)^{\frac{1}{2}} \\
\alpha=\left(\frac{\mu \sigma \omega}{2}\right)^{\frac{1}{2}} \\
\alpha=\sqrt{\frac{\mu \sigma \omega}{2}}
\end{gathered}
$$

Skin depth,

$$
\delta=\frac{1}{\alpha}
$$

$$
\delta=\sqrt{\frac{2}{\mu \sigma \omega}}
$$

Since $\omega=2 \pi f$, where $f$ is a linear frequency

$$
\delta=\sqrt{\frac{1}{\mu \sigma \pi f}}
$$

$>$ It is clear from the above expression that skin depth for good conductor depends upon the frequency of EM waves.

## For poor conductors or insulators:

$$
\frac{\sigma}{\omega \varepsilon} \ll 1
$$

Therefore

$$
\begin{gathered}
\left(1+\frac{\sigma^{2}}{\omega^{2} \varepsilon^{2}}\right)^{\frac{1}{2}} \approx 1+\frac{\sigma^{2}}{2 \omega^{2} \varepsilon^{2}} \\
\alpha=\omega\left\{\frac{\mu \varepsilon}{2}\left(1+\frac{\sigma^{2}}{2 \omega^{2} \varepsilon^{2}}-1\right)\right\}^{\frac{1}{2}} \\
\alpha=\omega\left(\frac{\mu \varepsilon}{2} \frac{\sigma^{2}}{2 \omega^{2} \varepsilon^{2}}\right)^{\frac{1}{2}}
\end{gathered}
$$

$$
\begin{aligned}
& \alpha=\frac{\sigma}{2}\left(\frac{\mu}{\varepsilon}\right)^{\frac{1}{2}} \\
& \alpha=\frac{\sigma}{2} \sqrt{\frac{\mu}{\varepsilon}} \\
& \delta=\frac{2}{\sigma} \sqrt{\frac{\varepsilon}{\mu}}
\end{aligned}
$$

$>$ From this expression it is clear that skin depth for insulator does not depend upon frequency of EM waves.

For a conducting medium, $\sigma=5.8 \times 10^{6} \mathrm{~S} / \mathrm{m}$ and $\varepsilon_{\mathrm{r}}=1$. Find conduction and displacement current densities if magnitude of electric field intensity $E$ is given by $E=150 \sin \left(10^{10} t\right) V / m .[2018-19]$

Conduction current density is given by $J_{c}=\sigma E$

$$
\begin{aligned}
\Rightarrow \quad J_{c} & =5.8 \times 10^{6} \times 150 \sin \left(10^{10} t\right) \mathrm{A} / \mathrm{m}^{2} \\
& =8.7 \times 10^{8} \sin \left(10^{10} t\right) \frac{A}{m^{2}}
\end{aligned}
$$

And,

$$
\begin{gathered}
\text { And, } J_{d}=\frac{\partial D}{\partial t}=\varepsilon \frac{\partial E}{\partial t} \\
\Rightarrow \quad J_{d}=\varepsilon_{r} \varepsilon_{0} \frac{\partial}{\partial t}\left\{150 \sin \left(10^{10} t\right)\right\} \\
\\
J_{d}=\quad J_{d}=1 \times 8.854 \times 10^{-12} \times 150 \times 10^{10} \times \cos \left(10^{10} t\right) \\
13.28 \cos \left(10^{10} t\right) \mathrm{A} / \mathrm{m}^{2}
\end{gathered}
$$

Calculate the magnitude of Poynting vector at the surface of the sun. Given the power radiated by sun is $5.4 \times 10^{\mathbf{2 8}}$ Watts and radius of the sun is $\mathbf{7 \times 1 0 ^ { 8 }} \mathbf{~ m}$.

We know Poynting vector is given by $S=\frac{P}{4 \pi r^{2}}$

$$
\begin{array}{ll}
\Rightarrow & S=\frac{3.8 \times 10^{26} \mathrm{~W}}{4 \times 3.14 \times\left(7 \times 10^{8} \mathrm{~m}\right)^{2}} \\
\Rightarrow & S=\frac{5.4 \times 10^{28} \mathrm{Joule}}{4 \times 3.14 \times\left(7 \times 10^{8}\right)^{2} \mathrm{~m}^{2} . \mathrm{S}} \\
\Rightarrow & S=8.77 \times 10^{9} \text { Joule } /\left(\mathrm{m}^{2} \cdot \mathrm{~s}\right)
\end{array}
$$

If the magnitude of $\overrightarrow{\boldsymbol{H}}$ in a plane wave is $1 \mathrm{Amp} /$ meter, find the magnitude of $\vec{E}$ for plane wave in free space. [2015-16]

$$
\begin{aligned}
& \text { We know that, } \frac{H_{o}}{E_{o}}=\sqrt{\frac{\varepsilon_{o}}{\mu_{o}}} \\
& \text { or } E_{o}=\sqrt{\frac{\mu_{o}}{\varepsilon_{o}}} H_{o} \\
& \qquad \begin{aligned}
E_{o} & =1 \times \sqrt{\frac{4 \pi \times 10^{-7}}{8.85 \times 10^{-12}}} \\
& =376.72 \mathrm{~V} / \mathrm{m}
\end{aligned}
\end{aligned}
$$

## Lecture 17

## ELECTROMAGNETIC FIELD THEORY



## CONTENTS

$>$ Numericals

Earth receives solar energy from the Sun which is $\frac{2 c a l}{\mathrm{~cm}^{2} . M i n}$. what are the amplitude of electric and magnetic fields of radiation?

The sunlight strikes upper atmosphere of earth with energy 1400 Watt $/ \mathrm{m}^{2}$. Calculate peak values of electric and magnetic fields at points.

$$
\begin{aligned}
& S=2 \mathrm{cal} /\left(\mathrm{cm}^{2} \cdot \mathrm{Min}\right) \\
& =\frac{2 \times 4.2 \mathrm{Joule}}{\left(10^{-2} \mathrm{~m}\right)^{2} \times 60 \mathrm{~s}} \\
& =1400 \text { Joule } /\left(\mathrm{m}^{2} \cdot \mathrm{~s}\right)
\end{aligned}
$$

We know Poynting vector is given by, $S=E X H=\mathrm{EH} \sin 90=\mathrm{EH}$

Therefore

$$
\begin{equation*}
E H=1400 \frac{\mathrm{Joule}}{\mathrm{~m}^{2} . \mathrm{s}} \ldots \tag{1}
\end{equation*}
$$

We also know by characteristic impedance of vacuum is given by

$$
\begin{equation*}
\frac{E}{H}=376.77 \Omega \tag{2}
\end{equation*}
$$

Multiplying equations (1) and equation (2) we get

$$
\begin{gathered}
E^{2}=527478(N / C)^{2} \\
\Rightarrow E=726.27 \mathrm{~N} / C
\end{gathered}
$$

Putting the value of $E$ in equation (1) we get,

$$
\begin{aligned}
H= & \frac{1400}{726.27} N /(A . m) \\
& H=1.93 N /(\text { A.m) }
\end{aligned}
$$

$$
\begin{gathered}
E_{0}=E \sqrt{2} \\
\\
E_{0}=726.27 \sqrt{2} \mathrm{~N} / \mathrm{C} \\
\Rightarrow \quad E_{0}=1027.1 \mathrm{~N} / \mathrm{C}
\end{gathered}
$$

$$
H_{0}=H \sqrt{2}
$$

$$
\Rightarrow
$$

$$
H_{0}=1.93 \sqrt{2} N /(A . m)
$$

$$
H_{0}=2.73 \mathrm{~N} /(\text { A. } \mathrm{m})
$$

A 100 Watt sodium lamp radiating its power. Calculate electric and magnetic fields strength at a distance of $5 \mathbf{m}$ from the lamp.

We know Poynting vector is given by, $\quad S=\frac{P}{4 \pi r^{2}}$

$$
\begin{array}{cc}
\Rightarrow & S=\frac{100 \mathrm{~W}}{4 \times 3.14 \times(5 \mathrm{~m})^{2}} \\
\Rightarrow & S=\frac{100 \mathrm{Joule}}{4 \times 3.14 \times 25 \mathrm{~m}^{2} \cdot \mathrm{~S}} \\
\Rightarrow & S=3.185 \times 10^{-1} \mathrm{Joule} /\left(\mathrm{m}^{2} \cdot \mathrm{~s}\right)
\end{array}
$$

We know $\quad S=\vec{E} X \vec{H}=\mathrm{EH} \sin 90=\mathrm{EH}$
Therefore

$$
\begin{equation*}
E H=3.185 \times 10^{-1} \mathrm{~J} /\left(\mathrm{m}^{2} . \mathrm{s}\right) \tag{1}
\end{equation*}
$$

We also know by characteristic impedance of vacuum is given by

$$
\begin{equation*}
\frac{E}{H}=376.6 \Omega \tag{2}
\end{equation*}
$$

Multiplying equations (1) and equation (2) we get

$$
\begin{gathered}
E^{2}=119.95(\mathrm{~V} / \mathrm{m})^{2} \\
E=10.95 \mathrm{~V} / \mathrm{m}
\end{gathered}
$$

Putting the value of E in equation (2), we get

$$
\begin{aligned}
& H=\frac{10.95}{376.6} \frac{\mathrm{amp}-\mathrm{turn}}{\mathrm{~m}} \\
& H=0.0291 \frac{\mathrm{amp}-\mathrm{turn}}{\mathrm{~m}}
\end{aligned}
$$

The sunlight strikes the upper atmosphere of earth with energy flux $1.38 \mathbf{k W m}^{-2}$. What will be the peak values of electric and magnetic fields at the points.
[2019-20]

As we know that ' $S$ ' is given by energy per unit time per unit area.

$$
\begin{gathered}
S=1.38 \mathrm{kWm}{ }^{-2} \\
S=1.38 \times 10^{3} \text { Joule } / \mathrm{m}^{2} \mathrm{~s}
\end{gathered}
$$

We know Poynting vector is given by, $S=E X H=\mathrm{EH} \sin 90=\mathrm{EH}$
Therefore

$$
\begin{equation*}
E H=1.38 \times 10^{3} \text { Joule } / m^{2} s \tag{1}
\end{equation*}
$$

We also know by characteristic impedance of vacuum is given by

$$
\frac{E}{H}=376.6 \Omega \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \text {................ }
$$

Multiplying equations (1) and equation (2) we get

$$
\begin{gathered}
E^{2}=519708(\mathrm{~V} / \mathrm{m})^{2} \\
\Rightarrow E=720.91 \frac{\mathrm{~V}}{\mathrm{~m}}
\end{gathered}
$$

Now,

$$
E_{0}=E \sqrt{2}=720.91 \times 1.414 \mathrm{~V} / \mathrm{m}
$$

$$
\Rightarrow \quad E_{0}=1.02 \mathrm{kV} / \mathrm{m}
$$

Putting the value of E in equation (1), we get ,

$$
H=\frac{\mathbf{1} .38 \times 10^{3}}{720.91} \mathrm{amp}-\mathrm{turn} / \mathrm{m}
$$

$\mathrm{H}=1.91 \mathrm{amp}-\mathrm{turn} / \mathrm{m}$
Now,

$$
\begin{aligned}
& H_{0}=H \sqrt{2} \\
\Rightarrow & \\
& H_{0}=1.91 \sqrt{2} \mathrm{amp}-\mathrm{turn} / \mathrm{m} \\
\Rightarrow \quad & H_{0}=2.7 \mathrm{amp}-\mathrm{turn} / \mathrm{m}
\end{aligned}
$$

For silver, $\mu=\mu_{0}$ and $\sigma=3 \times 10^{7} \mathrm{mhos} / \mathrm{m}$. Calculate the skin depth at $10^{8} \mathrm{~Hz}$ frequency. Given, $\mu_{0}=4 \pi \times 10^{-7} N / A^{2}$.
[2016-17]
Sol: Skin depth is given by ,

$$
\begin{gathered}
\delta=\sqrt{\frac{2}{\mu \sigma \omega}} \\
\delta=\sqrt{\frac{2}{\mu \sigma 2 \pi f}} \\
\delta=\sqrt{\frac{1}{\mu \sigma \pi f}}
\end{gathered}
$$

$$
\delta=\sqrt{\frac{1}{4 \pi \times 10^{-7} N / A^{2} \times 3 \times 10^{7} m h o / m \times \pi \times 10^{8} \mathrm{~Hz}}}
$$

$$
\delta=9.18 \times 10^{-6} \mathrm{~m}
$$

Using Maxwell's equation, curl $\vec{B}=\mu_{0}\left(\vec{J}+\frac{\partial \vec{D}}{\partial t}\right)$, prove that div $\vec{D}=\rho$
According to the given problem, curl $\vec{B}=\mu_{0}\left(\vec{J}+\frac{\partial \vec{D}}{\partial t}\right)$
Taking divergence of both sides, we get

$$
\begin{gathered}
\vec{\nabla} \cdot(\vec{\nabla} \times \vec{B})=\mu_{0} \vec{\nabla} \cdot\left(\vec{J}+\frac{\partial \vec{D}}{\partial t}\right) \\
\vec{\nabla} \cdot(\vec{\nabla} \times \vec{B})=0 \\
\vec{\nabla} \cdot\left(\vec{J}+\frac{\partial \vec{D}}{\partial t}\right)=0
\end{gathered}
$$

$$
\vec{\nabla} \cdot \vec{J}+\frac{\partial}{\partial t}(\vec{\nabla} \cdot \vec{D})=0
$$

From equation of continuity, $\vec{\nabla} \cdot \vec{J}=-\frac{\partial \rho}{\partial t}$
Hence, $\frac{\partial}{\partial t}(\vec{\nabla} \cdot \vec{D})=\frac{\partial \rho}{\partial t}$

$$
\vec{\nabla} \cdot \vec{D}=\rho
$$

Thank

